

Thermodynamic Stability and Anion Ordering of Perovskite Oxynitrides

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Abstract

Perovskite oxynitrides (PONs) are a promising class of materials for applications ranging from catalysis to photovoltaics. However, the vast space of single PON materials ($\text{ABO}_{3-x}\text{N}_x$) has yet to be fully explored. Additionally, the community needs guidelines that relate PON chemistry and anion ordering to stability to better understand how to design PON materials that resist corrosion and decomposition under operating conditions. Screening this materials space requires identifying candidate PON materials that will be stable under operating conditions, which in turn requires methods to evaluate each material's stability. Here we predict the stability of single PON materials using a four-step approach based on density functional theory modeling: (i) enumerate viable cation pairs, (ii) select an energetically favorable prototypical anion ordering, (iii) compute each PON's energy above the thermodynamic convex hull, and (iv) generate computational Pourbaix diagrams to determine allowable ranges of electrochemical operating conditions. A critical part of our approach is determining a prototypical stable anion ordering for both ABO_2N and ABON_2 stoichiometries across a variety of A- and B-site cations. We demonstrate a stable anion ordering containing a high degree of *cis* ordering between B cations and minority-composition anions. We predict 85 stable and 109 metastable PON compounds, with $A \in \{\text{La}, \text{Pb}, \text{Nd}, \text{Sr}, \text{Ba}, \text{Ca}\}$ and $B \in \{\text{Re}, \text{Os}, \text{Nb}, \text{Ta}\}$ forming cation pairs that lead to stable PONs less than 10 meV/atom above the thermodynamic convex hull. Computational Pourbaix diagrams for two stable candidates, CaReO_2N and LaTaON_2 , suggest that not all compounds with zero energy above the thermodynamic convex hull can be easily synthesized.

1 Introduction

Heteroanionic materials are a class of compounds in which cations (often transition metals, but possibly any metal, inorganic, or organic cation) coordinate with two or more anions. This class includes materials such as metal oxysulfides, metal oxynitrides, and metal oxyfluorides. Heteroanionic materials have several applications, including in energy storage as battery cathode materials,¹ in solid-state illumination as phosphor materials,² and in ammonia synthesis as electrocatalysts.³ The heteroanionic material design space is vast owing to the many possible combinations of metal cations, anion choice, anion

ordering, and lattice structures.¹ Perovskites are a subset of heteroanionic compounds characterized by at least two metal cations surrounded by anions arranged in polyhedral geometry. Perovskite materials are well known for their applications in photovoltaics,⁴ but have also been used for water splitting⁵ and light-emitting diodes.⁶

Perovskite oxynitrides (PONs) are a special case of perovskite oxides in which the anion sites are occupied by either O or N, with the general single-PON formula $\text{ABO}_{3-x}\text{N}_x$, with $0 < x < 3$. While some PONs with anion defects may have anions arranged on non-octahedral polyhedra that share faces or edges, in this work we consider only defectless perovskite structures with corner-sharing octahedra. Single PONs consist of 12-fold-coordinated A-site cations and 6-fold-coordinated B-site cations, with the B-site cations coordinated by octahedra of O or N anions that may arrange to form anion disorder, short-range order, or long-range order. Along with anion ordering and composition, the choice of A and B metal cations with multiple possible oxidation states creates a large combinatorial space for tuning the physical and electronic properties of these materials. At least 68 single PON materials have been synthesized, and many more potentially stable PONs have been predicted.⁷

Because the vast space of PON compounds has not yet been thoroughly explored, many promising material chemistries are likely unrealized. Understanding how chemical composition and anion ordering affect the stability of single PONs is critical to realize their synthesis and utility in several applications. Here we employ a density-functional-theory (DFT)-based, hierarchical down-selection approach to study PON stability across a wide compositional space of single PONs. We recognize that a full study of PONs would include the significantly larger design space of double and layered PONs, but for reasons of computational tractability we limit the scope of this work to the high tunability already possible in single PONs. This high compositional tunability may enable fine-tuning the physicochemical properties of these materials, making them useful for a variety of applications including electrochemical reactions,^{3,8} photocatalysis,⁹ photovoltaics,¹⁰ and ionic semiconduction.¹¹ **Figure 1** shows our four-part approach for finding stable single PONs. To create a feasible search space, we begin by enumerating cation pairs likely to form stable PON struc-

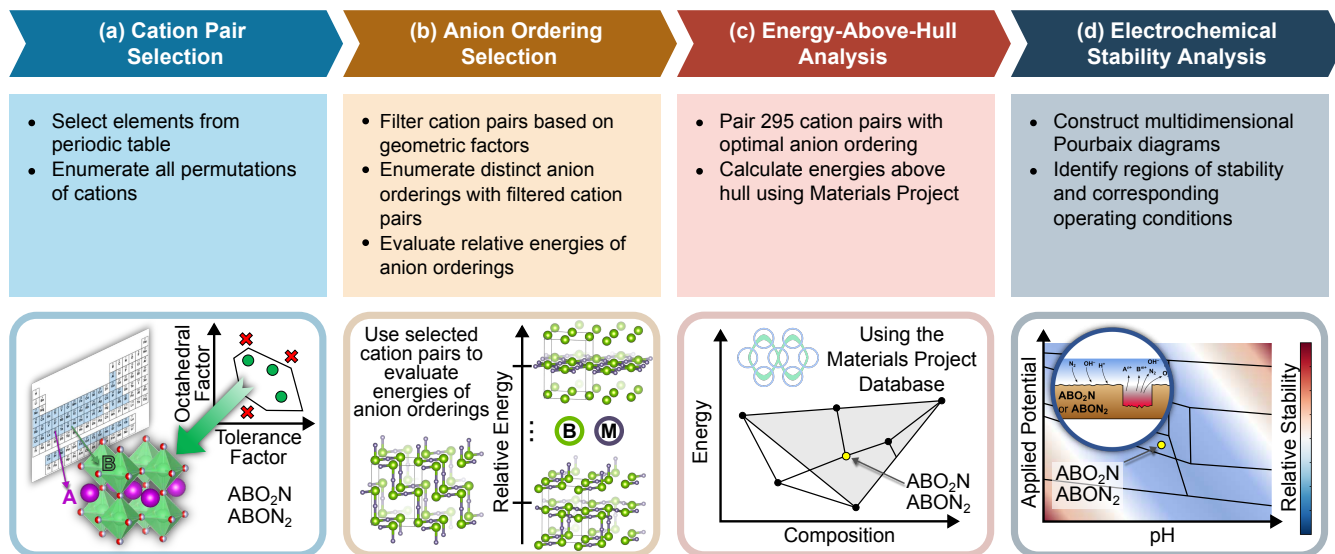


Figure 1. Workflow for identifying stable perovskite oxynitride compounds. The workflow involves four major steps: (a) cation pair selection, (b) anion ordering selection, (c) energy-above hull analysis, and (d) electrochemical stability analysis via computational Pourbaix diagrams.

tures using selected elements along with filtering rules to keep the number of DFT calculations feasible (Figure 1a). Second, we address the question of anion ordering by enumerating symmetrically distinct orderings of O^{2-} and N^{3-} within the single PON unit cell and evaluating their relative ground-state energies when paired with representative cation pairs (Figure 1b). Third, we combine the enumerated cation pairs and prototypical stable anion ordering determined in the first two steps to form a set of single PON compounds. For each single PON compound, we calculate the energy above the thermodynamic convex hull as a metric of stability (Figure 1c).¹² Finally, we explore the stability of two selected PONs, $CaReO_2N$ and $LaTaON_2$, in aqueous electrochemical reactions by creating Pourbaix diagrams and mapping regions of electrochemical stability as a function of applied potential and pH (Figure 1d). Our workflow demonstrates that *cis* ordering of O and N anions around the $B(O,N)_6$ octahedra is preferred above *trans* ordering, consistent with literature reports of synthesized perovskite oxynitrides and oxysulfides.^{13,14} Our screening identifies a set of 85 stable (≤ 10 meV/atom) and 109 metastable (> 10 meV/atom and ≤ 25 meV/atom) PONs, potentially tripling the space of known stable PON materials. The most stable compositions involve La, Pb, Nd, Sr, Ba, or Ca as the A-site cation and Re, Os, Nb, or Ta as the B-site cation. Evaluating the stability of $CaReO_2N$ and $LaTaON_2$ shows that while both are stable with regard to decomposition to other solid compounds, only $LaTaON_2$ is also stable under reasonable aqueous electrochemical operating conditions.

2 Methods

All DFT calculations were performed using the Vienna Ab Initio Simulation Package (VASP), version 5.4.4.^{15–18} Calculation preparation and post-processing were done using the Atomic Simulation Environment¹⁹ (version 3.17) and Pymatgen²⁰ (version 2020.1.28) libraries. Unless otherwise noted, all calculations used the same VASP settings used in the Mate-

rials Project (as implemented in the MPRelaxSet class in Pymatgen) to compare results with the Materials Project database.²¹ The projector-augmented wave method and associated pseudopotentials were used to describe core electrons,^{22,23} and the PBE exchange-correlation functional was used.²⁴ Calculations were spin-polarized, with plane waves expanded in a basis set up to 520 eV. k points were determined automatically according to the MPRelaxSet standard with a reciprocal k point density of $64 k$ points/ \AA^{-3} in reciprocal space. For each PON structure, we adjusted the ground-state energy to include corrections to oxide energies and GGA+ U energies that are necessary to combine our ground-state energies with Materials Project data.^{25,26}

The PON structures were modeled using an idealized cubic crystal $Pm\bar{3}m$ $\sqrt{2} \times \sqrt{2} \times 2$ cell with 20 atoms total and lattice vectors $(a, b, c) = (5.40776 \text{ \AA}, 5.40776 \text{ \AA}, 6.64773 \text{ \AA})$ based on material mp-4019 of the Materials Project database.²¹ Although we used this cubic crystal as the template for all geometry relaxations, we allowed for each PON to relax out of cubic symmetry. For geometry relaxations used to calculate PON decomposition energies, we conducted relaxations of cell volume, cell shape, and ion position (i.e., ISIF = 3) at least two times to alleviate errors arising from Pulay stresses. Relaxations of cell volume and cell shape also allow for each PON to adopt symmetry and octahedral tilts other than those specified in our idealized geometry. For geometry relaxations used to determine the prototypical anion ordering, we conducted at least two relaxations of cell shape and volume (ISIF = 2), followed by a relaxation of ion positions (ISIF = 1). Geometry relaxations proceeded using conjugate gradient or quasi-Newton force optimizers. For bulk relaxations of PON structures, we used a tighter electronic tolerance (10^{-5} eV) compared to the Materials Project default (10^{-3} eV) to help some structures converge. Where necessary, VASP mixing parameters were adjusted and RMM-DIIS iterations were replaced with blocked Davidson iterations to ensure self-consistent electronic convergence. Anion orderings were obtained for the 12 anion

sites in our $\sqrt{2} \times \sqrt{2} \times 2$ PON supercell. We enumerated anion orderings assuming fixed O_2N or ON_2 anion compositions, with each anion allowed to occupy any of the octahedral vertex sites around the B-site cations. The Enumlib software package was used to enumerate symmetrically distinct anion orderings under constraints of anion composition.²⁷

Five-dimensional computational Pourbaix diagrams were created for the Ca–Re–O–N–H and La–Ta–O–N–H systems. These many-component Pourbaix diagrams were constructed from the half-space intersection for all phases in each system, with the phase of lowest Pourbaix potential identified as the stable phase at the given conditions.^{28,29} Ion energies are obtained from the Materials Project Pourbaix diagram framework.³⁰ Details of the Pourbaix potential derivation are provided in Section S5.

3 Results and Discussion

The four parts of our workflow in Figure 1 filter the space of single PONs for stable candidates. Section 3.1 discusses our enumeration of cation pairs, including classification of in-hull or out-of-hull status based on the Goldschmidt tolerance and octahedral factors. Section 3.2 details this study’s principal effort to improve the accuracy of stability predictions by identifying a prototypical anion ordering that is generally preferred across all cation chemistries. Section 3.3 reports our screening process to identify stable PON structures by combining the cation pairs identified in Section 3.1 and the anion ordering identified in Section 3.2. Section 3.4 discusses using Pourbaix diagrams to further evaluate stability in the context of electrochemical reactions and highlights the need to supplementing general catalyst searches with application-specific screening.

3.1 Cation Pair Selection

As a first line of screening and to more clearly focus on the influence of cation chemistry and anion stoichiometry on stability, we focused on two PON compositions: ABO_2N and $ABON_2$. We selected A- and B-site cations from a set of 57 elements focusing on non-radioactive transition and post-transition metals, rare-earth metals, alkali metals, and alkaline metals (Figure S1). Allowable cation pairs were determined by enumerating pairs of A and B cations with nominal charges adding to +7 or +8 (i.e., to charge-balance the anions in ABO_2N or $ABON_2$ compositions, respectively), and for which A and B were not the same element. Pairs of cations with multiple oxidation states, such that given cation pair could add to both +7 and +8 charges, or to the same +7 or +8 charge in different ways, were counted as separate pairs. This enumeration process resulted in a set of 310 unique cation pairs. We note that a recent study by Wang et al.⁷ investigating the stability of oxynitrides, oxyfluorides, and nitrofluorides included a few Tc-containing oxynitride cation pairs not found by our enumeration process, specifically $Ba^{II}Tc^V O_2N$, $Ca^{II}Tc^V O_2N$, $Pb^{II}Tc^V O_2N$, $La^{III}Tc^V O_2N$, $Sr^{II}Tc^V O_2N$, and $La^{III}Tc^{IV} O_2N$. While our enumeration excludes radioactive elements, we include these six cation pairs for completeness and for comparison to that study, bringing the total number of unique cation pairs to 316.

Screening studies of perovskite structures often exploit geometric descriptors, such as the Goldschmidt tolerance³² and octahedral³³ factors based on ionic radii, to predict whether a metal oxide of certain cation chemistry will crystallize in perovskite geometry. Some PONs are known to be stable, meaning that experiments have verified their synthesizability into some type of perovskite geometry rather than another that of a non-perovskite polymorph. Proposed PON structures with Goldschmidt tolerance and octahedral factors close to those of known stable PONs are themselves more likely to prefer

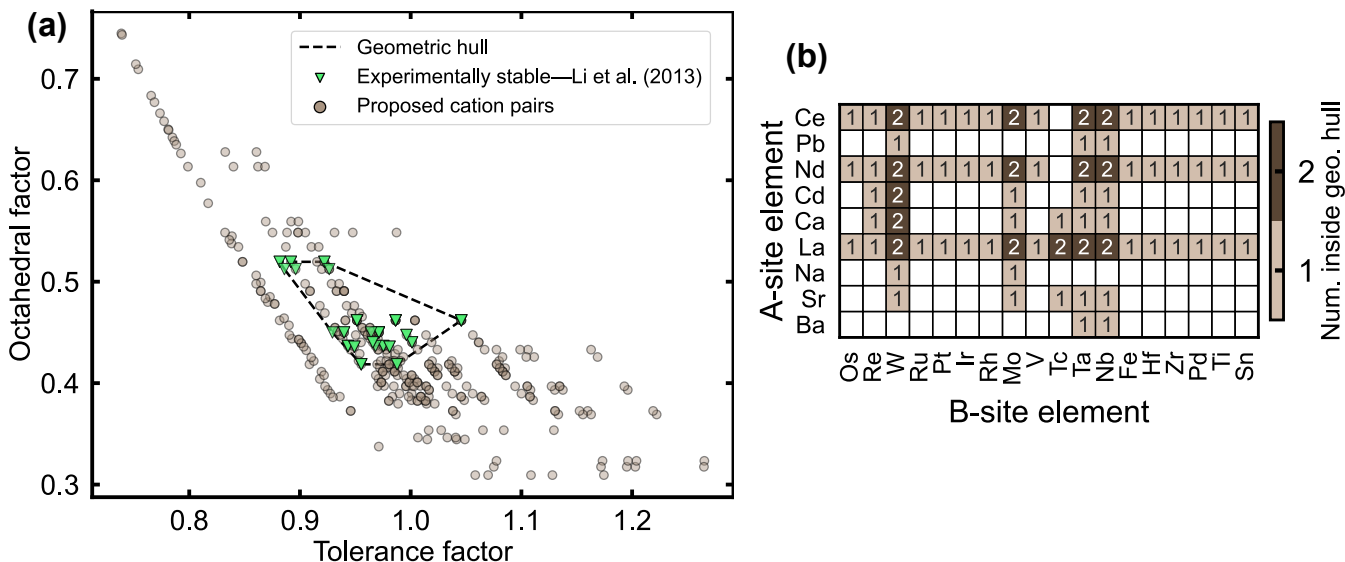


Figure 2. Perovskite oxynitride cation pairs predicted to fall within the geometric hull. (a) Structure map of Goldschmidt tolerance and octahedral factors for various single PON cation pairs. Experimentally stable points (green triangles) refer to PONs that have been synthesized and form an geometric hull (black dashed line).³¹ Proposed cation pairs (brown circles) refer to the PONs that we screened for stability. (b) Heatmap showing the counts of proposed PON compounds that fall inside the geometric hull. Blank cells indicate that no proposed PON of that cation pair falls within the geometric hull. All cation pairs with a count of 1 are ABO_2N .

perovskite geometry when crystallizing. For each of the 316 unique proposed cation pairs and the 68 cation pairs of known stable PONs, we calculated the Goldschmidt tolerance³² and octahedral³⁴ factors using geometric means of the ionic radii in each structure. **Table 1** shows how these geometric means are used to adapt the Goldschmidt tolerance and octahedral factors for perovskite oxynitrides. Here, r_A , r_B , r_O , and r_N are the ionic radii of the A-site cation, the B-site cation, an oxide anion, and a nitride anion, respectively. Oxidation-state-dependent Shan-

Table 1. Geometric-averaged Goldschmidt tolerance and octahedral factors for ABO_2N and $ABON_2$ perovskite oxynitrides.³¹

	Goldschmidt tolerance factor	Octahedral factor
ABO_2N	$\frac{[(r_A + r_O)^8(r_A + r_N)^4]^{1/12}}{\sqrt{2}[(r_B + r_O)^4(r_B + r_N)^2]^{1/6}}$	$\frac{r_B}{(r_O^4 r_N^4)^{1/6}}$
$ABON_2$	$\frac{[(r_A + r_O)^4(r_A + r_N)^8]^{1/12}}{\sqrt{2}[(r_B + r_O)^2(r_B + r_N)^4]^{1/6}}$	$\frac{r_B}{(r_O^2 r_N^4)^{1/6}}$

non ionic radii were obtained from Shannon’s 1976 revised tables.³⁵ The list of all cation pairs considered, with their Goldschmidt tolerance and octahedral factors, is given in **Table S1**.

The resulting data appear in **Figure 2a**, which displays a structure map for the proposed and known stable cation pairs. The 68 known stable cation pairs were used to form a convex hull shape against which we compare the 316 unique proposed cation pairs. To avoid confusion with the concept of a thermodynamic convex hull, we refer to the convex hull shape as the geometric hull. We consider a proposed PON with a certain cation pair likely to form perovskite geometry if its point on the structure map falls inside the geometric hull. Many proposed cation pair points cluster nearby or inside of the geometric hull, indicating that there may be many yet-undiscovered PONs that are experimentally stable. Of the 316 unique cation pairs, 90 (28.5%) fall within the geometric hull and 226 (71.5%) fall outside. Additionally, a number of points lie outside of but visually close to the geometric hull (i.e., within a perpendicular Euclidean distance of 0.02 units from an edge of the geometric hull). 28 of the 316 unique cation pairs fulfill this criterion. We later explore the possibility that the true region of stability may be different than that enclosed by the hull of currently synthesized materials. The 90 points inside the geometric hull mostly represent compounds with $A \in \{\text{Ce, Nd, La}\}$ and $B \in \{\text{Re, W, Mo, Ta, Nb}\}$ (**Figure 2b**). In particular, many compounds with $A = \text{La}$ or $B = \text{W}$ are likely to form in perovskite geometry, according to the Goldschmidt tolerance and octahedral factors. Our first-pass analysis using only geometric factors identifies potential cation chemistries that lead to stable PONs. However, the many points outside of but close to the geometric hull suggest that some stable PONs might remain undiscovered using this method. This finding motivates a deeper evaluation of stability with consideration of a prototypical anion ordering.

3.2 Determination of prototypical anion ordering

We next address the question of which anion orderings are most stable in single PONs. Screening PONs is complicated by the possibility of many anion orderings and the potential for each unique cation pair to prefer a different anion ordering, mak-

ing a DFT study of all possible combinations infeasible. Thus, the present study focuses on understanding how the preferred anion ordering changes as a function of cation charge configuration and identifying reasonable prototypical anion orderings that can be used for many structures at once. For heteroanionic perovskite systems of two anions in which the anions have the same charge but vastly different sizes, strain effects largely determine which anion orderings are preferred.¹⁴ This is the expected behavior for perovskite oxysulfides, in which O^{2-} and S^{2-} are predicted to have effective Shannon ionic radii of 1.40 Å and 1.84 Å, respectively.³⁵ However, if the two anions have similar sizes but very different charges, then the preferred anion orderings are likely those that minimize the total electrostatic energy of the structure and will thus depend on the charges of the cations. We expect this latter behavior with PON systems, in which the respective Shannon ionic radii of O^{2-} and N^{3-} are 1.40 Å and 1.46 Å.

We propose that the dependence of anion order on cation choice can be modeled by ranking the ground-state energies of a series of PON structures, each with the same cation pair but different anion orderings. While it would be computationally infeasible to calculate the ground-state energies of all possible anion orderings with all 316 unique cation pairs, we can feasibly test all anion orderings combined with a limited set of cation pairs. We selected 16 unique cation pairs to use for anion ordering ranking, focusing on a subset of cation pairs that both (i) spans a range of charge configurations and (ii) spans a range of A and B cation radii. The selected charge configurations were $A^I B^{VI}$, $A^{II} B^V$, and $A^{III} B^{IV}$ for $ABON_2$; and $A^I B^{VI}$ and $A^{III} B^V$ for ABO_2N . These criteria consider the effects both of cation charge (i.e., electrostatics) and cation size (i.e., strain effects) that might affect the preferred anion ordering. **Table S3** shows the list of cation pairs used for anion ordering analysis, along with their charges and ionic radii. We preferred cation pairs in which at least one of the cations had no more than one nominal oxidation state. The rationale behind this preference is that limiting the number of oxidation states reduces the number of cation pairs that could be matched with both ABO_2N and $ABON_2$ compositions, thus helping clearly distinguish the effect of PON composition on the preferred anion ordering.

To find possible anion orderings, we began with an ideal $ABO_3 \sqrt{2} \times \sqrt{2} \times 2$ perovskite structure and explored partial substitution of O^{2-} with N^{3-} anions, subject to the constraints that O^{2-} and N^{3-} compositions remain in a 2:1 ratio for ABO_2N or 1:2 ratio for $ABON_2$. We identify 32 symmetrically distinct anion orderings satisfying these constraints. These anion orderings are shown in **Figure S2** for ABO_2N and **Figure S3** for $ABON_2$. Because of the symmetry of the ABO_2N and $ABON_2$ compositions, the 32 anion orderings are the same between compositions except for the swapping of O^{2-} and N^{3-} anions.

We next combined each of the 16 unique cation pairs chosen for this analysis with each of the 32 anion orderings, choosing either the ABO_2N anion orderings for cation pairs adding to a charge of +7 or the $ABON_2$ anion orderings for cation pairs adding to a charge of +8. In total, we evaluated the ground-state energies of 512 structures. We calculate the ground-state energies using DFT rather than point-ion electrostatic models such as the Madelung energy³⁶ for reasons of accuracy (**Figure S7**). These calculations represent about 5% of the calculations needed to exhaustively determine the prototypical

anion ordering for all 316 unique cation pairs discussed earlier. For each cation pair, the anion orderings were ranked in order of increasing energy (Figures S8–S12 and Tables S4–S8), with the lowest-energy anion ordering for each cation pair assigned a ranking number of 1, the second-lowest a ranking number of 2, and so on. Cumulative rankings were calculated across all cation pairs tested by summing these ranking numbers and by summing the relative ground-state energies, leading to the prediction of one anion ordering in particular as most preferred on average over all cation pairs (Figure S13). We

henceforth denote this first most preferred anion ordering, both in Figure S13 and in the main text, as O_0 , with the second most preferred anion ordering denoted O_1 , and so on.

Figure 3 shows the O_0 anion ordering and key results of the anion ordering analysis. Figure 3a–d contains four different views of anion ordering O_0 . Figure 3a shows a full O_0 PON structure which has an ABO_2N composition and arbitrary A and B cations. Figure 3b shows a $2 \times 2 \times 4$ supercell with an alternative view of the O_0 anion ordering focusing on the topology of M–B–M bonds between the B-site cation and the minority-

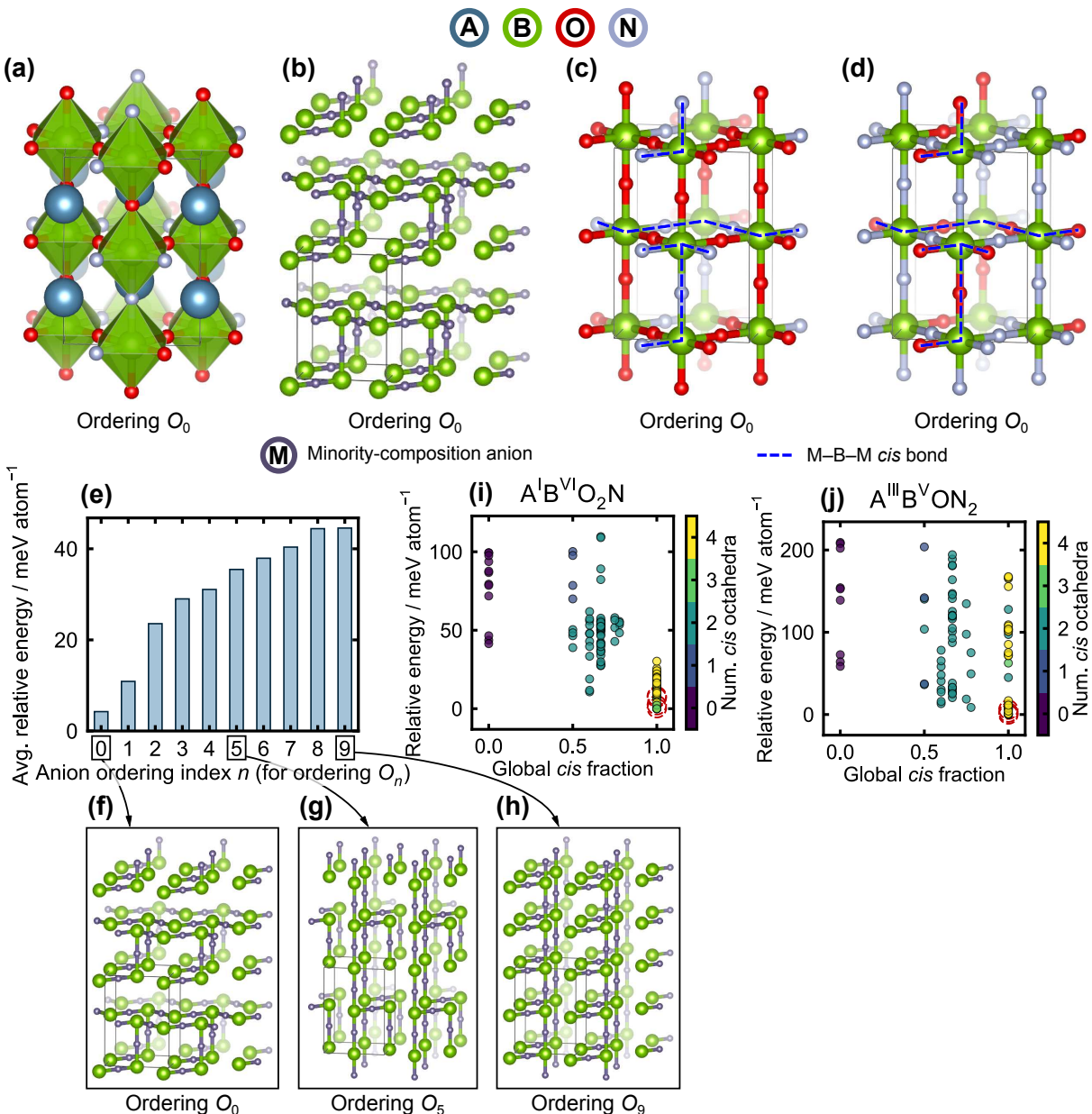


Figure 3. Rationalizing proposed anion ordering for general PON screening. All atom renderings are of ideal (unrelaxed) geometry, while all energies shown are for DFT-relaxed geometry. (a) $\sqrt{2} \times \sqrt{2} \times 2$ supercells showing the full structure of our proposed prototypical, stable anion ordering O_0 . (b) $2 \times 2 \times 4$ supercell showing the topology of the O_0 anion ordering with only the M–B–M bonds and no A-site cations. M (dark gray) is the minority-composition anion (i.e., N for ABO_2N and O for $ABON_2$). (c, d) $\sqrt{2} \times \sqrt{2} \times 2$ supercells showing the O_0 anion ordering, with A-site cations omitted for clarity, for (c) ABO_2N and (d) $ABON_2$ structures. Dotted lines overlay the M–B–M *cis* bonds. (e) Top 10 cumulative rankings by average ground-state energy for each anion ordering, across all 16 cation pairs tested and across all 32 symmetrically distinct anion orderings (Figure S13), with orderings O_0 , O_5 , and O_9 (f–h) shown for comparison. (i, j) Calculated relative DFT energy versus metrics of the amount of *cis* ordering for two representative cation charge configurations. Points surrounded by dotted red circles indicate the O_0 anion ordering (Figure S6).

composition (M) anion. **Figure 3c, d** shows the skeletal framework of B–(O, N) bonds for ABO_2N and ABON_2 compositions, respectively. These two structures are identical except for the swapping of O and N elements. **Figure 3e** shows the rankings of O_0 and a few other anion orderings as measured by relative DFT energy averaged across all 16 cation pairs used for anion ordering analysis (see **Figure S13** for rankings of all orderings). This ranking predicts O_0 to be thermodynamically preferred over all other anion orderings, on average. **Figure 3f–h** shows the $2 \times 2 \times 4$ supercells of bonding topology for a few selected anion orderings of different averaged relative DFT energy. Finally, **Figure 3i, j** shows how certain measures of the degree of M–B–M *cis* bonding relate to ground-state energy.

In PON structures where the compositions of each anion are not equal, changing the positions of M–B–M bonds will affect the degree of M–B–M *cis* bonding in the structure. The topology and degree of M–B–M *cis* bonding are important factors in rationalizing why certain anion orderings are preferred in PON structures. The degree of *cis* ordering within each B-site octahedron has been shown to correlate to thermodynamic stability in other heteroanionic perovskite compounds. For example, a recent computational study that exhaustively evaluated all possible anion orderings of the metal oxysulfide $\text{SrHf}(\text{O}_{0.5}\text{S}_{0.5})_3$ predicted that the preferred anion ordering must have anions arranged in a *cis* configuration around the B-site cation.¹⁴ This preferred ordering was rationalized in terms of competing electronic, electrostatic, and strain effects. Additionally, a neutron and electron-diffraction study of SrBO_2N (B = Nb, Ta) PONs found that SrBO_2N consistently prefers a partial anion ordering with –B–M– bonds forming *cis*-oriented chains within the equatorial planes of the BO_2N octahedra.¹³ The present work tests the hypothesis that similar *cis* orderings are preferred for perovskite oxynitrides generally.

The preferred O_0 anion ordering contains many instances of *cis* bonding. To rationalize quantitatively the importance of *cis* bonding in PON compounds (and why some *cis*-ordered anion orderings are lower in energy than others), we calculated two simple metrics of the degree of *cis* bonding: the global fraction of *cis* bonding and the number of *cis* octahedra. For each anion ordering, the global fraction of *cis* bonding is calculated by counting all of the M–B–M bonds within a $\sqrt{2} \times \sqrt{2} \times 2$ supercell and determining which fraction of those bonds have a 90° angle through the B-site cation. The number of *cis* octahedra is calculated by counting how many of the four octahedra in a $\sqrt{2} \times \sqrt{2} \times 2$ supercell contain at least one *cis* M–B–M bond (**Figures S4** and **S5**). These metrics were chosen to measure not just how often M–B–M *cis* bonding occurs in an anion ordering, but also whether anion orderings with M–B–M *cis* bonds distributed throughout the structure are preferred about those with such bonds concentrated in one part of the structure. **Figure 3i, j** shows these metrics for all 32 anion orderings when combined with cation pairs of charge configurations $\text{A}^{\text{III}}\text{B}^{\text{V}}$ and $\text{A}^{\text{I}}\text{B}^{\text{VI}}$, respectively, with data for the remaining charge configurations shown in **Figure S6**. Although there is too much variance in the relative ground-state energies of each anion ordering to state strong quantitative correlations, some trends exist. The points near global *cis* fractions of 0.0, 0.6, and 1.0 form three general clusters. The point at the bottom of each cluster indicates the anion ordering leading to the minimum possible ground-state energy for a certain global *cis*

fraction, similar to a thermodynamic convex hull plot. As the global *cis* fraction increases, the bottom point of each cluster becomes lower in energy, indicating that a higher fraction of *cis* bonds enables lower-energy anion orderings. A similar relationship exists between the relative ground-state energy and the number of *cis* octahedra. The clusters at a 0.0 global *cis* fraction tend to be composed of anion orderings with *cis* bonds distributed across zero or one octahedra. The clusters at a 0.6 global *cis* fraction generally have anion orderings with *cis* bonds distributed across two or three octahedra, and for a 1.0 global *cis* fraction, *cis* bonds are usually distributed across all octahedra in the anion ordering. In general, anion orderings with zero or one *cis* octahedra are higher in ground-state energy than those with three or four *cis* octahedra. The anion orderings that are first-, second-, and third-lowest in ground-state energy each have three *cis* octahedra. The most important observation of **Figure 3i, j** is that for each of the 16 cation pairs used for anion ordering analysis, the lowest-energy anion ordering is always one that has a 100% global fraction of *cis* bonding in M–B–M bonds. No anion ordering with even a single *trans* M–B–M bond surpasses the ground-state stability of an anion ordering with all *cis* M–B–M bonds, even if those *cis* bonds are unevenly distributed among the octahedra.

To restate, we used a heuristic strategy to identify a prototypical anion ordering preferred broadly across a wide range of cation chemistries. We enumerated all 32 possible symmetrically distinct orderings of O and N anions within $\sqrt{2} \times \sqrt{2} \times 2$ supercell. We then evaluated the energy of each anion ordering when separately combined with 16 different cation pairs representative of typical cation charge configurations and chemistries. Based on the averaged ground-state energies of the anion orderings across all 16 cation pairs, we predict anion ordering O_0 as an acceptable prototype for the preferred anion ordering across all cation chemistries. Bond-counting statistics suggest that the prototypical anion ordering must have M–B–M bonds which are always in a *cis* configuration and never in a *trans* configuration. Although other factors such as the octahedral tilt and the particular atomic decoration of each octahedron may influence the energy of each anion ordering and deserve further investigation, the major finding is that a highly *cis* configuration correlates generally with low energy.

3.3 Analysis of Energy Above Hull

Thermodynamic stability can be directly estimated in the absence of experimental data using quantum-mechanical methods. Such calculations can predict the energy above the thermodynamic convex hull, or the thermodynamic driving force (i.e., decomposition energy) for a PON material to decompose to its elements, metal nitride or oxide counterparts, or other products.³⁷ With 316 unique cation pairs and a prototypical anion ordering identified, we compute the the energy above the convex hull to determine PON structure stability. On a plot of ground-state energy versus material composition, the convex hull is the locus of points representing polymorphs or structures with the lowest energy for a given composition. Materials on or close to the convex hull are less likely to decompose to other products.

We computed the energy above the convex hull for each of our PON materials. Our original set of 316 cation pairs in-

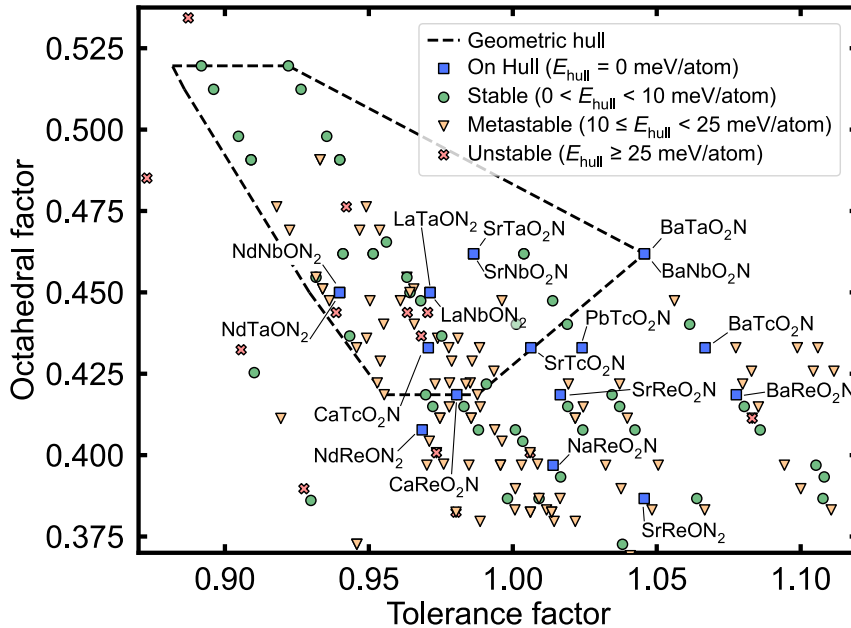


Figure 4. Predicted stability classes for selected PON structures inside or near the geometric hull. Blue squares refer to structures on the thermodynamic convex hull, according to our DFT calculations. Green circles, yellow triangles, and red X marks respectively refer to stable, metastable, and unstable structures, based on decomposition energy ranges as defined in the text. The dotted black line is the geometric hull from Figure 2, based on data from Li et al.³¹

cludes the possibility for some cation pairs to have multiple charge configurations due to elements with multiple possible oxidation states. However, plane-wave DFT computations usually consider just the composition of a PON structure without the particular oxidation state of the cations, so our set of 316 unique cation pairs reduces to 295 unique cation compositions. For each of the 295 cation compositions, the cations were combined with the O_0 anion ordering to form a full PON structure. These 295 PON structures are then subjected to bulk geometry relaxation to determine the bulk ground-state energy per atom for each structure. Of these relaxations, 227 converged to perovskite geometry under two or more consecutive combined volume-ionic relaxations. The ground-state energies of these 227 converged structures were then combined with the Materials Project database to calculate each structure’s decomposition energy relative to the lowest-energy linear combination of stoichiometrically equivalent decomposition products.

Based on the calculated energies above hull, we classify each material’s stability from a DFT and thermodynamic standpoint. Figure 4 shows the results of energy-above-hull calculations for a selected subset of cation pairs that are inside or near the geometric hull from Figure 2. Table 2 shows the stability classifications of the geometric hull (Section 3.1) versus the stability classifications obtained from DFT-based energy-above-hull calculations (Section 3.3). Thermodynamic computations place the median energy above hull for metastable inorganic crystalline materials at 15 meV/atom.¹² However, for the subset of nitride compounds, this median is 63 meV/atom, owing to the remarkable stability of the N_2 molecule and its strong triple bond. Because our PON structures contain only A–N, B–N, and O–N nitride bonds, even if they are not triple bonds, we reason that some of these materials may still possess metastability at energies above hull higher than the typical inorganic

crystalline thresholds. To focus our study on those compounds most likely to be synthesizable, we adopt conservative cutoffs for defining stable, metastable, and unstable PONs. For an energy above hull E_{hull} , we define $0 < E_{\text{hull}} < 10$ meV/atom as stable, $10 \text{ meV/atom} \leq E_{\text{hull}} < 25$ meV/atom as metastable, and $E_{\text{hull}} \geq 25$ meV/atom as unstable. $E_{\text{hull}} = 0$ meV/atom signifies materials that are on the thermodynamic convex hull. Figure 4 shows that a large number of stable materials and materials on the thermodynamic convex hull fall within an approximate $0.95 \leq t \leq 1.07$ and $0.375 \leq \mu \leq 0.475$ window, where t and μ are the Goldschmidt tolerance and octahedral factors, respectively. This window extends outside of the geometric hull, indicating that there may be several potentially promising materials not captured by the geometric hull. Among the compounds predicted to have zero energy above hull are a few known to be synthesizable, including $SrNbO_2N$ and $LaTaON_2$ ^{38–40} as well as $LaNbON_2$, $BaNbO_2N$, $NdNbON_2$, $NdTaON_2$, $SrTaO_2N$, and $BaTaO_2N$.³¹

DFT-derived thermodynamic design rules offer critical additional insight beyond that of the simpler geometric hull. Of the 227 PON structures converged under DFT calculations for which we obtained thermodynamic energies above the convex hull, 85 are classified stable or on-hull, 33 are classified unstable, and the rest are classified metastable (see Table 2). For each of the 118 stable, on-hull, and unstable materials, we determined whether the geometric hull predicts that the material will form perovskite geometry (i.e., whether the material is plotted inside the geometric hull in Figure 2). While our DFT-calculated stability classes do not necessarily represent ground truth, the mismatch of geometric hull predictions with DFT-predicted stability classes warrants concern. For example, the geometric hull predicts perovskite geometry for stable/on-hull PONs and non-perovskite geometry for unstable PONs for only

Table 2. Counts of PONs by geometric hull classification (“Inside geo. hull” and “Outside geo. hull”) versus by stability class (“On Hull”, “Stable”, “Metastable”, and “Unstable”) derived from energy-above-hull calculations on DFT ground-state energies. “On Hull” refers to a compound having zero energy above the thermodynamic convex hull, not to being on the boundary of the geometric hull.

Stability Class	Num. inside geo. hull	Num. outside geo. hull	All
On Hull ($E_{\text{hull}} = 0$ meV/atom)	9	11	20
Stable ($0 < E_{\text{hull}} < 10$ meV/atom)	27	38	65
Metastable ($10 \leq E_{\text{hull}} < 25$ meV/atom)	37	72	109
Unstable ($E_{\text{hull}} \geq 25$ meV/atom)	7	26	33
All	80	147	227

52.5% of materials. For individual stability classes, the geometric hull predicts perovskite geometry for 45.0% of on-hull materials and 41.5% of stable materials, and non-perovskite geometry for 78.8% of unstable materials. Overall, the geometric hull’s prediction of whether a cation chemistry prefers perovskite geometry over other geometry correlates poorly with the calculated thermodynamic stability of that cation chemistry when constrained to a perovskite oxynitride structure. We conclude that in a search for stable PON compounds, the geometric hull approach can initially help exclude unstable candidate materials, but will likely fail to identify many potentially stable materials. Calculating the energy above the thermodynamic convex hull more clearly describes the propensity for a PON solid to remain stable and the degree to which it resists decomposition into thermodynamically competing material phases than if we were to rely on simpler, more common structure rules such as the Goldschmidt tolerance and octahedral factors.

Figure 5 shows the decomposition enthalpy of each perovskite oxynitride structure with a specific cation pair and the anion ordering for its stoichiometry. Several trends follow from this plot. First, there are more stable or metastable ABO_2N compounds than ABON_2 compounds. Of the 227 compounds converged under DFT relaxation, 156 are of stoichiometry ABO_2N and 71 are of stoichiometry ABON_2 (see Tables S9 and S10). The lower count of ABON_2 compounds arises from the fact that relatively fewer cation pairs both have geometrically preferred sizes and sum to a +8 cationic charge. While 42.3% of the ABO_2N compounds are classified as stable, only 26.8% of ABON_2 compounds meet this threshold. The trend is reversed for metastable compounds, with 43.0% of ABO_2N and 59.2% of ABON_2 compounds ranked metastable. However, unstable compounds make up approximately equal percentages of each stoichiometry class (14.7% and 14.1% for ABO_2N and ABON_2 , respectively).

Second, perovskite stability generally increases as the size of the B-site cation decreases. Many more perovskite oxynitride structures classified as stable or metastable appear on the left sides of both heatmaps than the right sides, corresponding with smaller B-site cations. A recent computational study of perovskite oxides identified a similar trend, predicting that the probability of perovskite oxide formability, even across multiple cation charge configurations, increases when the A-site cation is sufficiently larger than the B-site cation.⁴¹ This observation is consistent with the distribution of points in Figures 2 and 4, in which the bulk of stable PON compounds and PON compounds with zero energy above hull are concentrated at higher values of the Goldschmidt tolerance factor (i.e., where A has a higher ionic radius compared to B). Additionally, the

study also found the effect to be enhanced for A^{III} cation pairs over A^{I} and A^{II} cation pairs. Interestingly, our data show different trends, with stable or on-hull PONs dominated by A^{II} compounds (46.3% of all stable or on-hull), followed by A^{III} (36.3%), A^{I} compounds (12.5%), and finally A^{IV} compounds (5.0%) (Tables S11 and S12). Stability appears to be favored when the B-site cation has a more positive charge than the A-site cation does. This may be explained by the need for the B-site cation to coordinate six negatively charged anions while the 12 ions coordinating the A-site cation include a mix of cations and anions. A less positive A-site cation may also have a larger ionic radius, corresponding to the increased stability observed with relatively smaller B-site cations.

Third, perovskite oxynitride structures in which Re is the B-site cation are predicted to be stable, with energies above hull of less than 23 meV/atom across all A-site cations and across both perovskite oxynitride compositions considered. The present work predicts nearly all Re-containing compounds to be stable, in agreement with the conclusion of Wang et al.⁷ that Re-containing compounds represent a new class of potentially stable perovskite oxynitrides. A few other material classes, while not as stable as the $\text{B} = \text{Re}$ class of materials, are potentially promising. Materials with $\text{B} = \text{Os}$ are nearly all stable for ABO_2N stoichiometry, and nearly all metastable for ABON_2 stoichiometry. Similarly, many compounds with $\text{B} = \text{W}$ and $\text{B} = \text{Mo}$ for both stoichiometries are predicted to be metastable. These findings agree with a computational study of ternary metal nitrides, which predicted many potentially stable metal nitrides not found in the Inorganic Crystal Structure Database (ICSD).⁴² Among the ternary metal nitrides that study proposed for inclusion in the ICSD are compounds containing Re, Os, and Ru, some of which match in metal chemistry with the PONs proposed in this work. Of the stable and on-hull PONs in the present work, 8 with $\text{B} = \text{Re}$, 8 with $\text{B} = \text{Os}$, and 4 with $\text{B} = \text{Ru}$ match metal chemistries of proposed stable ternary metal nitrides. If we include metastable PONs, there are 15 with $\text{B} = \text{Os}$ and also 7 with $\text{B} = \text{W}$. Given that the limit of metastability is higher for nitride-containing materials than for inorganic crystalline materials generally, we hypothesize that even the metastable PONs predicted in this work may yet contain synthesizable PON compounds.

Additionally, there are many stable and metastable materials of both ABO_2N and ABON_2 stoichiometries with $\text{A} \in \{\text{Ce}, \text{Nd}, \text{Ca}, \text{La}, \text{Sr}, \text{Ba}\}$. In particular, we predict that many $\text{A} = \text{La}$ and $\text{A} = \text{Nd}$ compounds will be synthesizable, including B-site cations over a wide range of ionic radii.

As with the trends in B-site cation, some trends in A-site cation have connections to nitride materials. We note the recent synthesis of some perovskite nitride materials with similar

cation chemistries as those we predict to be stable in perovskite oxynitrides, specifically CeWN_{3-x} , CeMoN_{3-x} , and LaWN_3 perovskite nitrides.^{43,44} Synthesis of these nitrides often involves thin-film deposition via sputtering in an atmosphere with high levels of reactive nitrogen or in the presence of N plasma in order to avoid the incorporation of O into the nitride lattice. Nevertheless, nitride crystals are known to release N anions from the lattice, which facilitates both Mars-van Krevelen reaction mechanisms and the deactivation of nitride electrocatalysts.³ The propensity for pure nitrides to oxidize may make this work's perovskite oxynitride trends useful for synthesizing perovskite nitrides. Appropriate nitride synthesis techniques could be used to make stable or metastable perovskite nitrides with cation chemistry similar to the stable and metastable PONs in this work.

We note that each PON may have other competing oxynitride polymorphs that may be more stable. However, including these polymorphs in our energy-above-hull analysis could easily double or triple our study's computational cost as many of these structures do not exist in the Materials Project database

or similar databases. In addition, many of these oxynitride polymorphs will have some degree of anion site symmetry and will thus require additional calculations to consider possible anion orderings. We reiterate that our use of cell-volume and cell-shape relaxations already allows us to consider non-cubic perovskite polymorphs. Additionally, we omitted from energy-above-hull analysis any PON structure that relaxed out of perovskite geometry. In making these decisions, we have focused on the potential for a PON to decompose into oxide, nitride, or oxynitride crystals that already exist in the Materials Project database, even as we acknowledge that other, more stable polymorphs might exist.

Our energy-above-hull calculations have combined 295 unique cation chemistries with a prototypical optimal O_0 anion ordering. 227 of the resulting geometry relaxations converged and were used in combination with data from the Materials Project database to calculate the per-atom energy above the convex hull. In total, we find 85 stable compounds, 109 metastable compounds, and 33 unstable compounds. In particular, we find that cation pairs with $A \in \{\text{La, Pb, Nd, Sr, Ba, Ca}\}$

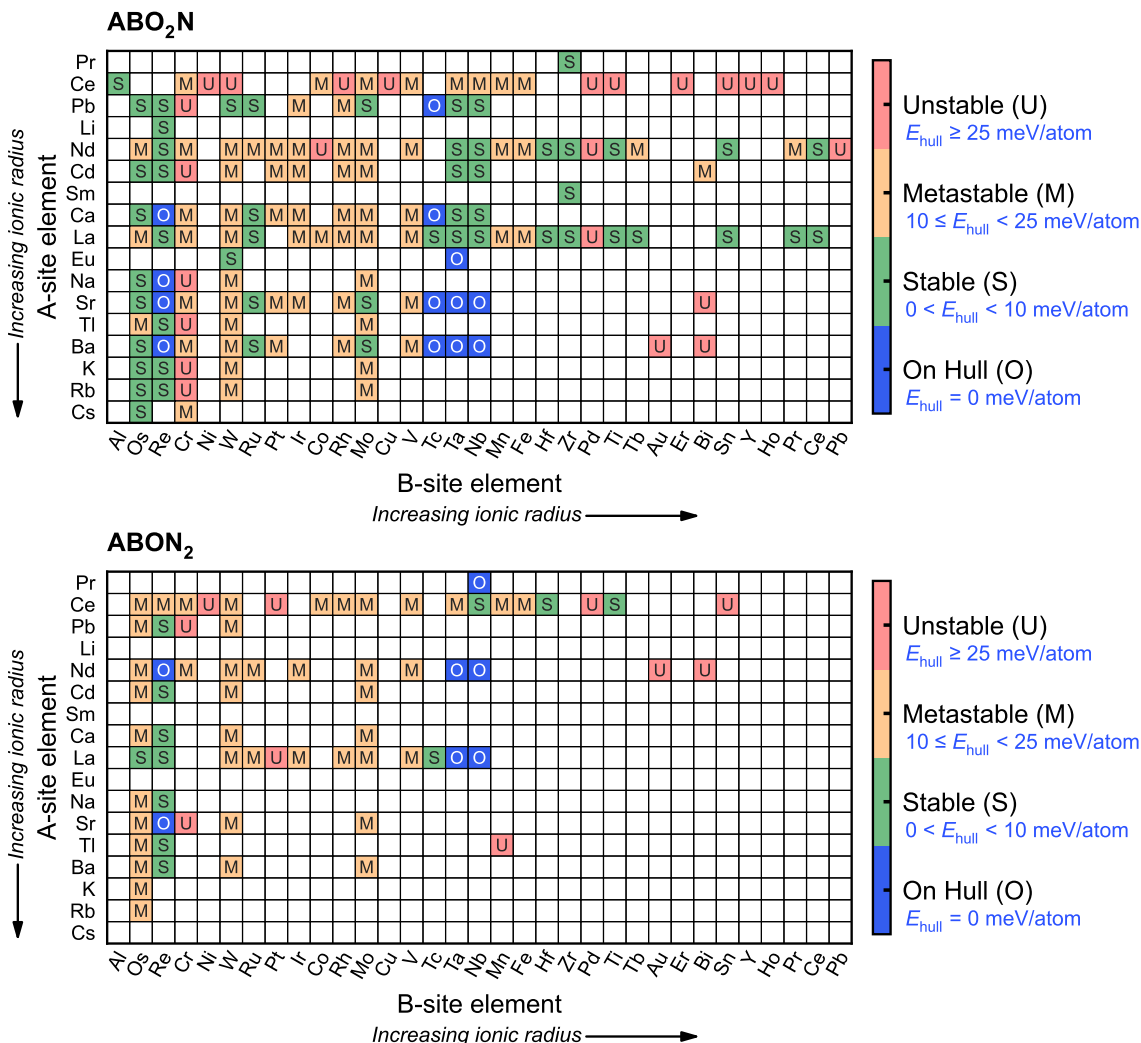


Figure 5. Heat map showing energy above hull for different cation pairs. Blue, green, yellow, and red, cells respectively indicate PON compounds classified as on-hull, stable, metastable, and unstable, based on decomposition energy ranges as defined in the text. White cells indicate cation pairs that are predicted to be highly unstable, not charge-balanced, or do not satisfy basic geometry requirements ($r_A > r_B$). The A- and B-site cations on each axis are arranged in order of increasing nominal ionic radius.³⁵

and $B \in \{\text{Re}, \text{Os}, \text{Nb}, \text{Ta}\}$ are likely to lead to stable or metastable PON compounds. Counting both metastable and stable PONs in our study as possibly synthesizable (totaling 194), our work may more than triple the number of PONs known to be experimentally synthesizable (i.e., compared to the 68 compounds in Li et al.³¹). The compounds identified could represent new heteroanionic materials for industrial, energy, and environmental materials applications.

3.4 Analysis of Electrochemical Stability

Perovskite oxide materials have been proposed for a number of important electrochemical reactions, such as the oxygen reduction, oxygen evolution, and hydrogen evolution reactions.^{45,46} Additionally, PON compounds are likely to catalyze surface reactions via the same types of electrochemical Mars-van Krevelen mechanisms observed in perovskite oxides for oxygen evolution⁴⁵ and in metal nitrides for ammonia synthesis.³ Thus, understanding a PON's electrochemical stability in addition to its thermodynamic stability is necessary to determine the PON's practical usability as an electrocatalyst. Pourbaix diagrams help predict whether a certain set of operating conditions—an operating pH, applied electrical potential—is capable of maintaining the PON in its solid state rather than promoting decomposition to competing aqueous, solid, or gaseous phases.⁴⁷ For example, the oxygen evolution reaction in acidic media often requires applied potentials of 2.0 V vs. RHE or higher for practical operation,⁴⁸ many transition metal nitrides (such as Mo, Fe, V, and Ni nitrides) are solid in acidic media only for potentials more negative than 0 V vs. SHE.⁴⁹ Transition metal nitrides may have more stability as acidic-media ammonia synthesis catalysts, where more negative applied potentials promote higher surface coverage of protons.³ Additionally, one may generate stability processing diagrams, which show regions of stability similar to a Pourbaix diagram, but in terms of two or more element activities instead of an electrical potential.^{50,51} Stability processing diagrams help identify the concentrations of species necessary to keep a phase stable. Beyond knowing a PON is thermodynamically stable outside of solution, Pourbaix and stability processing diagrams can help predict the conditions in which it will remain stable.

To this end, we compute multidimensional Pourbaix and stability processing diagrams to understand the operating conditions under which two candidate PON compounds would remain in a stable solid phase.^{29,52} Multicomponent Pourbaix diagrams have been used to study quaternary systems such as metal oxychalcogenides.⁵³ However, these diagrams are often constructed by solving the Nernst equation to draw phase boundaries between competing phases, an approach that cannot consider the free energy or concentration of each species individually. In contrast, our computational Pourbaix diagrams are derived using a thermodynamic grand potential for each species, which enables us to consider the energy contributions of both ion concentration and chemical potential for all elements.^{28,54} To the best of our knowledge, this work represents the first time this method has been applied to quaternary oxynitride systems.

We focus on two candidate PON compounds, CaReO_2N and LaTaON_2 , which were predicted in the previous step to have zero energy above the thermodynamic convex hull. **Figure 6**

shows Pourbaix diagrams for these two compounds. For such quaternary systems, the Pourbaix diagram has five dimensions: pH (linked to the chemical potential of O through water equilibrium), the redox potential E , and the chemical potentials of N, A-site cations, and B-site cations. The panels of **Figure 6** show two- or three-dimensional slices of this five-dimensional diagram. Note that standard Pourbaix diagrams are plotted at ion concentrations of 10^{-6} mol/L = 10^{-6} M, below the concentration at which a solid oxynitride might crystallize in solution. In this work, we express the A- and B-site chemical potentials as concentrations in **Figure 6a, c**, and plot **Figure 6b, d** at a higher ion concentration of 10^{-2} M. **Table S13** contains the calculated standard Gibbs free energy of formation for each phase in **Figure 6b, d**.

Figure 6a, c illustrates the ion concentrations and N chemical potentials at which the oxynitride remains solid, as well as other phases that compete at lower metal, O, and N chemical potentials. For both CaReO_2N and LaTaON_2 , stability is generally independent of the B-site ion concentration but depends strongly on $\Delta\mu_{\text{N}}$, the value of the N chemical potential relative to $\text{N}_2(\text{g})$. For both the Ca–Re and La–Ta systems, higher N chemical potentials favor the formation of a solid oxynitride phase over hydride or oxide species. In addition to decomposing to simpler metal oxides, nitrides, hydrides, or ions, we also consider the possibility that the PON could convert to its opposite composition (i.e., from CaReO_2N to CaReON_2 or from LaTaON_2 to LaTaO_2N). In **Figure 6a**, stability of the solid O_2N PON is favored at higher Ca and Re concentrations, but converts to its ON_2 form at low Ca concentrations. Additionally, CaReO_2N is the preferred phase only for a narrow window of N chemical potentials near $\Delta\mu_{\text{N}} \approx 4.0$ eV/atom. However, in **Figure 6c** and over the same concentrations of A- and B-site cations, LaTaON_2 is the only stable PON phase. The fact that both O_2N and ON_2 phases appear in **Figure 6a** but not in **Figure 6c** can be rationalized by the interplay of each phase's energy and sensitivity to the N chemical potential. CaReON_2 has an energy above hull of 0.40 meV/atom (i.e., higher in energy than CaReO_2N), but is also stabilized with increasing $\Delta\mu_{\text{N}}$ more than CaReO_2N due to its higher N content. Thus, there exists a small CaReO_2N region representing the locus of operating conditions for which the lower energy above hull of CaReO_2N outweighs the stabilization of N-rich CaReON_2 due to N chemical potential. In contrast, LaTaO_2N is a full 6.45 meV/atom higher in energy than LaTaON_2 , and there are no operating conditions for which LaTaO_2N becomes lower in energy than LaTaON_2 . Overall, **Figure 6a, c** suggests that the formability of a solid PON phase is controlled primarily by $\Delta\mu_{\text{N}}$, with possible influence from cation concentrations on whether that phase is an O_2N or an ON_2 phase.

Figure 6b, d shows that when a solid oxynitride phase is present, it forms generally only under reducing and alkaline conditions. In particular, even at $\Delta\mu_{\text{N}} = +2.5$ eV/atom vs. $\text{N}_2(\text{g})$, solid CaReON_2 would be formable only in very alkaline $\text{pH} \geq 12$ environments. In contrast, LaTaON_2 at $\Delta\mu_{\text{N}} = 0$ eV/atom has a much wider stability region that extends into acidic conditions, almost to $\text{pH} = 4$, but at the cost of requiring more negative applied potentials. Thus, LaTaON_2 might be synthesizable in alkaline, neutral, or slightly acidic environments.

Figure 6a shows that the formation of a Ca–Re PON is

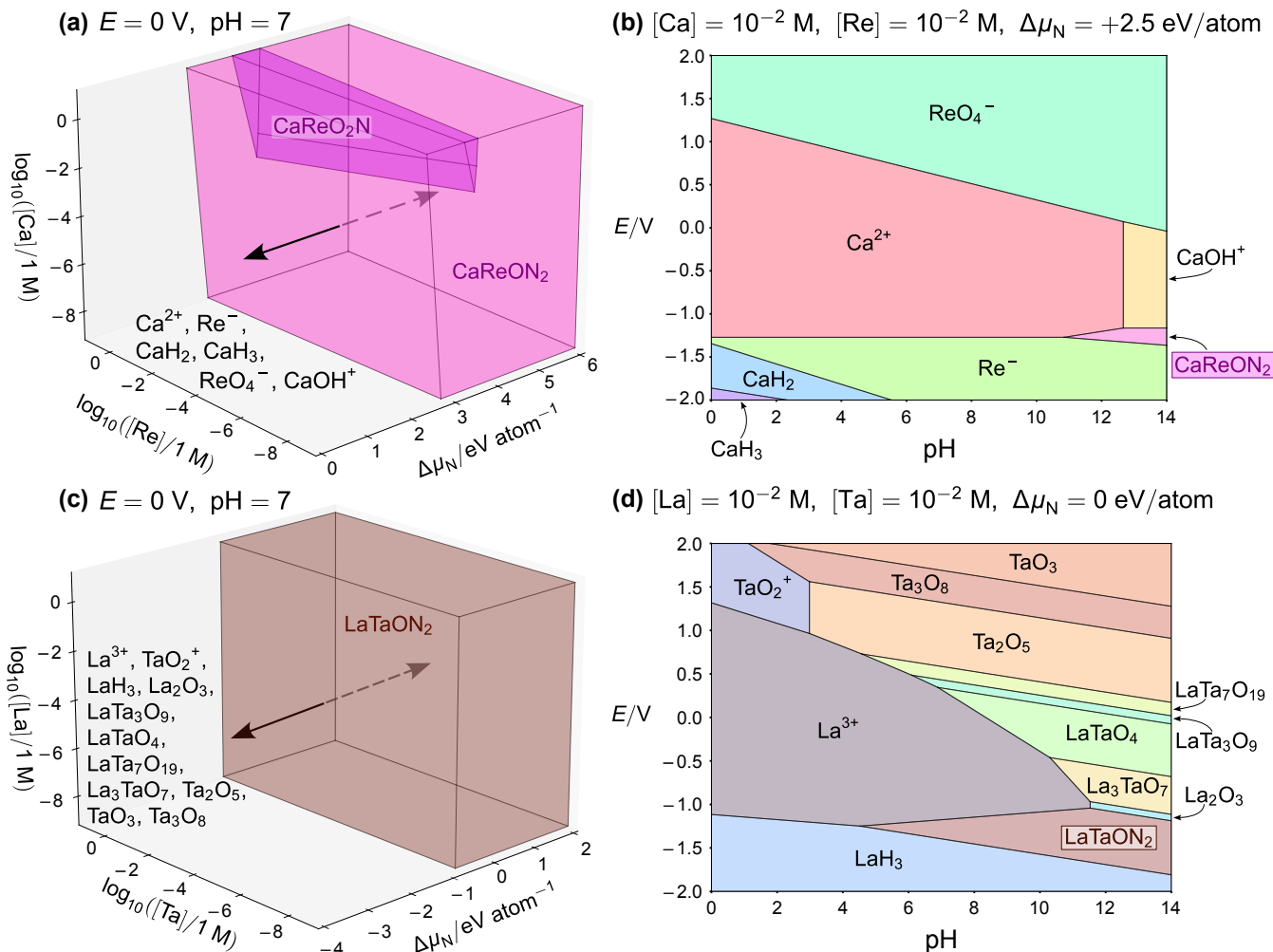


Figure 6. Computational stability processing and Pourbaix diagrams for CaReO_2N and LaTaON_2 . Selected projections of the 5D phase diagram for (a, b) CaReO_2N and (c, d) LaTaON_2 . (a, c) Stability processing diagrams in terms of A- and B-site ion concentration and N chemical potential. For simplicity, only regions corresponding to a solid oxynitride are shown, with other possible decomposition products written outside these regions. Arrows indicate changes in N chemical potential (with dashed arrows indicating the inside of a solid oxynitride region). (b, d) 2D Pourbaix projections of the 5D phase diagram, in terms of E and pH. Potentials are given in V vs. SHE. The boxed compound indicates a solid oxynitride phase. In all panels, neutral compounds are solid and ionic species are aqueous.

likely impossible in water, as CaReON_2 can be stabilized only for $\Delta\mu_{\text{N}} \gtrsim 2.5$ eV/atom (and even higher $\Delta\mu_{\text{N}}$ for CaReO_2N). Note that even plasma-cracked atomic nitrogen yields a N chemical potential of only $\Delta\mu_{\text{N}} \approx 1$ eV/atom,⁵⁵ which means it is highly improbable that an aqueous nitrogen precursor can attain the necessary N chemical potential required to form solid CaReO_2N . However, solid LaTaON_2 might be synthesized in an aqueous environment, as **Figure 6c** indicates stability of the solid PON at a N chemical potential of only $\Delta\mu_{\text{N}} = -1$ eV/atom, suggesting that the solid oxynitride might be stable in the aqueous state using nitrogen-rich precursors such as urea,⁵⁶ ammonia,⁵⁷ hydrazine,⁵⁸ or melamine.^{59,60} Indeed, multiple experimental studies report successful laboratory synthesis of LaTaON_2 or LaTaO_2N and confirm that this oxynitride does crystallize in a perovskite structure.^{38,61} Despite the lack of established benchmarking for their aqueous chemical potentials, such precursors have demonstrated efficacy in the synthesis of other nitrides and oxynitrides,^{59,62} suggesting their potential utility as precursors for the synthesis of PONs.

Our electrochemical analysis shows that a solid PON compound having zero energy above the thermodynamic hull according to DFT calculations may still require special operating conditions to be synthesizable or remain stable in aqueous solution. The energy above hull is useful for quickly screening many PON candidate compounds. However, further stability analysis via ab initio thermodynamics and experimental studies is needed to confirm the synthesizability and long-term stability of a given PON compound under a certain set of operating conditions. Tools such as many-component Pourbaix diagrams can focus high-throughput screening studies on just those candidates that are possible to synthesize at all.

4 Conclusion

We computationally investigate the thermodynamic stability and anion ordering of 156 ABO_2N and 71 ABON_2 perovskite oxynitride compounds using density functional theory calculations and Materials Project data.²¹ We identify a prototypical

anion ordering that is lowest in energy on average across 16 cation chemistries and five charge configurations. This prototypical anion ordering contains exclusively *cis* bonds connecting the B-site cation and the minority-composition anions, similar to that of metal oxysulfides.¹⁴

Based on our screening calculations, we predict 85 PON compounds to be stable (i.e., energies above hull of 10 meV/atom or less), of which 8 have already been demonstrated stable under laboratory conditions.³¹ These 85 stable compounds include PONs with $A \in \{\text{La, Pb, Nd, Sr, Ba, Ca}\}$ and $B \in \{\text{Re, Os, Nb, Ta}\}$. We predict that these PONs, in addition to the class of Re-containing compounds,⁷ will also be experimentally stable.

We generate computational Pourbaix diagrams for the CaReO_2N and LaTaON_2 stable candidate materials, which reveal that some materials with zero energy above hull still differ significantly in the range of chemical potential conditions that allow for the PON compound to remain solid. Specifically, a Ca–Re solid PON is likely not synthesizable under any practically accessible electrochemical conditions, characterized by an unphysically high required N chemical potential. In contrast, LaTaON_2 has a much wider region of operating conditions that could potentially allow for synthesizability of the solid LaTaON_2 PON in an alkaline, reducing environment. Ab initio thermodynamics and experiments are necessary to verify the stability of the proposed stable PONs in this study.

Our work has relevance beyond the question of whether a particular PON compound is stable. Many of the stable and metastable PONs identified in our study may be useful for electrocatalytic, photocatalytic, and photovoltaic applications. Additionally, the concept of enumerating classes of potentially stable PON materials by A- and B-site cation could be used to screen other compounds with similar compositional variability, such as double perovskites, layered perovskites, and spinels.

Supporting Information File 1: Detailed information on the enumeration of cation pairs within the geometric hull; additional information on the selection of preferred anion orderings; summary statistics of energy-above-hull calculations; and a derivation of the Pourbaix potential used to generate multi-component computational Pourbaix diagrams.

File 2: A zip file containing all 227 relaxed full PON structures, as well as all anion ordering structures, in both Pymatgen Structure JSON format and CIF format; and spreadsheets showing the raw DFT energy of each relaxed full PON structure and the relative energy of each anion ordering structure.

Author Contributions S.D.Y. wrote the original manuscript draft and performed all DFT calculations and analysis under the guidance of B.R.G. and G.P. J.C. performed all calculations for the Pourbaix and stability processing diagrams under the guidance of W.S. All authors helped write the manuscript. B.R.G. and G.P. conceived the original project idea and secured funding.

Notes The views expressed in this article do not necessarily represent the views of the Department of Energy or of the U.S. Government. The authors declare no competing financial interest.

Acknowledgements This work was supported in part by the U.S. Department of Energy through the Los Alamos National Laboratory. Los Alamos National Laboratory is operated by Triad National Security, LLC, for the National Nuclear Security Administration of the U.S. Department of Energy (contract no. 89233218CNA000001). Funding for this work was provided by the Laboratory Directed Research and Development program of Los Alamos National Laboratory under project no. 20200294ER and by startup funds of Bryan Goldsmith provided by the University of Michigan. This research also used resources of the National Energy Research Scientific Computing Center (NERSC), which is supported by the Office of Science of the U.S. Department of Energy under contract no. DE-AC02-05CH11231.

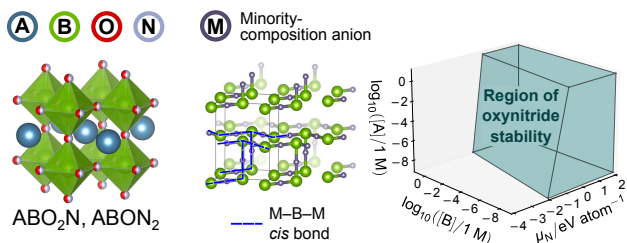
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A ∈ {La, Pb, Nd, Sr, Ba, Ca}
B ∈ {Re, Os, Nb, Ta}

Preferred ordering has 100% *cis* M-B-M bonding

Pourbaix diagrams predict conditions for aqueous stability

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